## Article

# On Discrete Fractional Solutions of Non-Fuchsian Differential Equations 

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#### Abstract

In this article, we obtain new fractional solutions of the general class of non-Fuchsian differential equations by using discrete fractional nabla operator $\nabla^{\eta}(0<\eta<1)$. This operator is applied to homogeneous and nonhomogeneous linear ordinary differential equations. Thus, we obtain new solutions in fractional forms by a newly developed method.


Keywords: discrete fractional calculus; fractional nabla operator; non-Fuchsian equations

## 1. Introduction

The history of fractional mathematics dates back to Leibniz (1695). This field of work is rapidly increasing and, nowadays, it has many applications in science and engineering [1-4]. Heat transfer, diffusion and Schrödinger equation are some fields where fractional analysis is used.

A similar theory was started for discrete fractional analysis and the definition and properties of fractional sums and differences theory were developed. Many articles related to this topic have appeared lately [5-18].

In 1956 [5], differences of fractional order was first introduced by Kuttner. Difference of fractional order has attracted more interest in recent years.

Diaz and Osler [6], defined the notion of fractional difference as follows

$$
\Delta^{\varsigma} \Phi(t)=\sum_{k=0}^{\infty}(-1)^{k}\binom{\varsigma}{k} \Phi(t+\varsigma-k)
$$

where $\varsigma$ is any real number.
Granger and Joyeux [19] and Hosking [20], defined notion of the fractional difference as follows

$$
\begin{aligned}
\nabla^{\varsigma} \Phi(t) & =(1-q)^{\varsigma} \Phi(t) \sum_{k=0}^{\infty}(-1)^{k} \frac{\Gamma(\varsigma+1)}{\Gamma(k+1) \Gamma(\varsigma-k+1)} q^{k} \Phi(t) \\
& =\sum_{k=0}^{\infty}(-1)\binom{\varsigma}{k} \Phi(t-k)
\end{aligned}
$$

where $\varsigma$ is any real number and $q \Phi(t)=\Phi(t-1)$ is the shift operator. Gray and Zhang [21], Acar and Atici [10] studied on a new definition and characteristics of the fractional difference.

## 2. Preliminary and Properties

In this section, we first present sufficient fundamental definitions and formulas so that the article is self-contained.

The rising factorial power $t^{\bar{m}}(t$ to the $m$ rising, $m \in \mathbb{N}$ ) is defined by

$$
t^{\bar{m}}=t(t+1)(t+2) \ldots(t+m-1), t^{\overline{0}}=1
$$

Let $\sigma$ be any real number. Then " $t$ to the $\sigma$ rising" is defined to be

$$
\begin{equation*}
t^{\bar{\sigma}}=\frac{\Gamma(t+\sigma)}{\Gamma(t)}, t \in \mathbb{R}-\{\ldots,-2,-1,0\}, 0^{\bar{\sigma}}=0 \tag{1}
\end{equation*}
$$

Also, the $\nabla$ operator of Equation (1) is given by

$$
\begin{align*}
& \nabla\left(t^{\bar{\sigma}}\right)=\nabla \frac{\Gamma(t+\sigma)}{\Gamma(t)} \\
= & \frac{\Gamma(t+\sigma)}{\Gamma(t)}-\frac{\Gamma(t-1+\sigma)}{\Gamma(t-1)} \\
= & \frac{(t-1+\sigma) \Gamma(t-1+\sigma)}{\Gamma(t)}-\frac{\Gamma(t-1+\sigma)}{\Gamma(t-1)}  \tag{2}\\
= & \frac{\Gamma(t-1+\sigma)}{\Gamma(t-1)}\left(\frac{t-1+\sigma}{t-1}-1\right) \\
= & \sigma \frac{\Gamma(t-1+\sigma)}{\Gamma(t)} \\
= & \sigma t^{\overline{\sigma-1}}
\end{align*}
$$

where $\nabla u(t)=u(t)-u(t-1)$.
Let $\eta \in \mathbb{R}^{+}$such that $m-1 \leq \eta<m, m \in \mathbb{N}$. The $\eta$ th-order fractional nabla sum of $g$ is given by

$$
\begin{equation*}
\nabla_{b}^{-\eta} g(t)=\frac{1}{\Gamma(\eta)} \sum_{s=b}^{t}(t-\delta(s))^{\overline{\eta-1}} g(s), \tag{3}
\end{equation*}
$$

where $t \in \mathbb{N}_{b}=\{b\}+\mathbb{N}_{0}=\{b, b+1, b+2, \ldots\}, b \in \mathbb{R}, \delta(s)=s-1$ is backward jump operator. Also, we define the trivial sum by $\nabla_{b}^{-0} g(t)=g(t)$ for $t \in \mathbb{N}_{b}$.

The $\eta$ th-order Riemann-Liouville type nabla fractional difference of $g$ is defined by

$$
\begin{align*}
\nabla_{b}^{\eta} g(t) & =\nabla^{m}\left[\nabla^{-(m-\eta)} g(t)\right]  \tag{4}\\
& =\nabla^{m}\left[\frac{1}{\Gamma(m-\eta)} \sum_{s=b}^{t}(t-\delta(s))^{\overline{m-\eta-1}} g(s)\right]
\end{align*}
$$

where $g: \mathbb{N}_{b}^{+} \longrightarrow \mathbb{R}[10]$.
Theorem 1 ([16]). Let $f$ and $g: \mathbb{N}_{0}^{+} \longrightarrow \mathbb{R}, \gamma, \phi>0$. Then

$$
\begin{gather*}
\nabla^{-\gamma} \nabla^{-\phi} f(t)=\nabla^{-(\gamma+\phi)} f(t)=\nabla^{-\phi} \nabla^{-\gamma} f(t),  \tag{5}\\
\nabla^{\gamma}[h f(t)+k g(t)]=h \nabla^{\gamma} f(t)+k \nabla^{\gamma} g(t), h, k \in \mathbb{R}  \tag{6}\\
\nabla \nabla^{-\gamma} f(t)=\nabla^{-(\gamma-1)} f(t),  \tag{7}\\
\nabla^{-\gamma} \nabla f(t)=\nabla^{(1-\gamma)} f(t)-\binom{t+\gamma-2}{t-1} f(0) . \tag{8}
\end{gather*}
$$

Lemma 1 (Power Rule [10]). Let $v>0$ and $\eta$ be two real numbers so that $\frac{\Gamma(\eta+1)}{\Gamma(\eta+v+1)}$ is defined. Then,

$$
\nabla_{b}^{-v}(t-b+1)^{\bar{\eta}}=\frac{\Gamma(\eta+1)}{\Gamma(\eta+v+1)}(t-b+1)^{\overline{\eta+v}}, \quad t \in \mathbb{N}_{b}
$$

Lemma 2 (Leibniz Rule [10]). For any $\eta>0$, $\eta$ th-order fractional difference of the product $f g$ is given in this form

$$
\begin{equation*}
\nabla_{0}^{\eta}(f g)(t)=\sum_{m=0}^{t}\binom{\eta}{m}\left[\nabla_{0}^{\eta-m} f(t-m)\right]\left[\nabla^{m} g(t)\right] \tag{9}
\end{equation*}
$$

where

$$
\binom{\eta}{m}=\frac{\Gamma(\eta+1)}{\Gamma(m+1) \Gamma(\eta-m+1)}
$$

and $f, g$ are defined on $\mathbb{N}_{0}$, and $t$ is a positive integer.
Lemma 3 (Index Law). Let $g(t)$ is single-valued and analytic. Then

$$
\begin{equation*}
\left(g_{\gamma}\right)_{\eta}=g_{\gamma+\eta}=\left(g_{\eta}\right)_{\gamma} \quad\left(g_{\gamma} \neq 0 ; g_{\eta} \neq 0 ; \gamma, \eta \in \mathbb{R} ; t \in \mathbb{C}\right) \tag{10}
\end{equation*}
$$

## 3. Main Results

We start by considering the following differential equation

$$
\begin{equation*}
\left(1+\frac{\ell}{x}\right) \frac{d^{2} y}{d x^{2}}+\left[a+\frac{b}{x}\left(1+\frac{\ell}{x}\right)\right] \frac{d y}{d x}+\left[c+\frac{d}{x}+\frac{\varepsilon}{x^{2}}\left(1+\frac{\ell}{x}\right)\right] y(x)=\psi \tag{11}
\end{equation*}
$$

where $\psi$ is a given function, $x \in \mathbb{C} \backslash\{0,-\ell\}$, and $a, b, c, d, \varepsilon$ and $\ell$ are parameters.
Let

$$
\begin{equation*}
y(x)=x^{\tau} e^{\kappa x} w(x) \tag{12}
\end{equation*}
$$

so that

$$
\begin{equation*}
\frac{d y}{d x}=x^{\tau-1} e^{\kappa x}\left[x \frac{d w}{d x}+(\tau+\kappa x) w(x)\right] \tag{13}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{d^{2} y}{d x^{2}}=x^{\tau-2} e^{\kappa x}\left[x^{2} \frac{d^{2} w}{d x^{2}}+2(\tau+\kappa x) x \frac{d w}{d x}+\left\{\kappa^{2} x^{2}+2 \tau \kappa x+\tau(\tau-1)\right\} w(x)\right] \tag{14}
\end{equation*}
$$

By substituting (12)-(14) into the (11), we have

$$
\begin{align*}
& x^{2}(x+\ell) \frac{d^{2} w}{d x^{2}}+\left[(2 \tau+b) \ell+(2 \tau+2 \kappa \ell+b) x+(2 \kappa+a) x^{2}\right] x \frac{d w}{d x} \\
& +[\{\tau(\tau+b-1)+\varepsilon\} \ell+\{\tau(\tau+2 \kappa \ell+b-1)+\kappa b \ell+\varepsilon\} x \\
& \left.+\left\{\kappa^{2} \ell+(2 \tau+b) \kappa+\tau a+d\right\} x^{2}+\left(\kappa^{2}+\kappa a+c\right) x^{3}\right] w(x) \\
& =x^{3-\tau} e^{-\kappa x} \psi(x), \quad x \in \mathbb{C} \backslash\{0,-\ell\} . \tag{15}
\end{align*}
$$

Finally, we find it to be suitable to restrict the different parameters involved in (11) and (15) by means of the following equalities;

$$
\begin{align*}
2 \tau+b & =0 \\
\tau(\tau+b-1)+\varepsilon & =0  \tag{16}\\
\kappa^{2}+\kappa a+c & =0
\end{align*}
$$

so that

$$
\begin{equation*}
\tau=-\frac{1}{2} b=\frac{-1 \pm \sqrt{1+4 \varepsilon}}{2} \tag{17}
\end{equation*}
$$

and

$$
\begin{equation*}
\kappa=\frac{-a \pm \sqrt{a^{2}-4 c}}{2} \tag{18}
\end{equation*}
$$

Under the parametric constraints given by (16), the Equation (15) will immediately decrease to a simpler form

$$
\begin{equation*}
(x+\ell) \frac{d^{2} w}{d x^{2}}+[2 \kappa \ell+(2 \kappa+a) x] \frac{d w}{d x}+\left(\kappa^{2} \ell+\tau a+d\right) w(x)=x^{1-\tau} e^{-\kappa x} \psi(x) \tag{19}
\end{equation*}
$$

where $\tau$ and $\kappa$ are given by (17) and (18), respectively.
Theorem 2. Let $w, \psi \in\left\{w, \psi: 0 \neq\left|w_{\eta}(x)\right|,\left|\psi_{\eta}(x)\right|<\infty\right\}$, and $\eta \in \mathbb{R}$. Then the nonhomogeneous linear differential equation

$$
\begin{equation*}
w_{2}(\alpha x+\beta)+w_{1}(\gamma x+v \alpha+\delta)+v \gamma w(x)=\psi(x), x \neq-\frac{\beta}{\alpha}, \alpha \neq 0, v \in \mathbb{R} \tag{20}
\end{equation*}
$$

has particular solutions in the below forms:

$$
\begin{align*}
& w^{I}(x)=\left\{\left[\psi_{-q^{-1} v}(\alpha x+\beta)^{\left(\delta \alpha-\gamma \beta-\alpha^{2}\right) / \alpha^{2}} e^{\frac{\gamma}{\alpha} x}\right]_{-1}(\alpha x+\beta)^{(\gamma \beta-\delta \alpha) / \alpha^{2}} e^{-\frac{\gamma}{\alpha} x}\right\}_{-1+q^{-1} v}  \tag{21}\\
& w^{I I}(x)=(\alpha x+\beta)^{\frac{-(\gamma x+\delta)}{\alpha}-v+1} \\
& \times\left(\left\{\left[\psi(\alpha x+\beta)^{\frac{\gamma x+\delta}{\alpha}+v-1}\right]_{q^{-1} v}(\alpha x+\beta)^{\frac{-\delta \alpha+\gamma \beta}{\alpha^{2}}+1} e^{-\frac{\gamma}{\alpha} x}\right\}_{-1}\right. \\
&\left.\times(\alpha x+\beta)^{\frac{\delta \alpha-\gamma \beta}{\alpha^{2}}-2} e^{\frac{\gamma}{\alpha} x}\right)_{-1-q^{-1} v} \tag{22}
\end{align*}
$$

where $w_{n}=\frac{d^{n} w}{d x^{n}}(n=0,1,2), w_{0}=w=w(x), \alpha, \beta, \gamma, v, \delta$ are given constants.
Proof. For $\psi(x) \neq 0$,
(i) When we operate $\nabla^{\eta}$ to the both sides of (20), we have

$$
\begin{equation*}
\nabla^{\eta}\left[w_{2}(\alpha x+\beta)\right]+\nabla^{\eta}\left[w_{1}(\gamma x+v \alpha+\delta)\right]+\nabla^{\eta}(w v \gamma)=\nabla^{\eta} \psi \tag{23}
\end{equation*}
$$

by using (9) and (10) we obtain

$$
\begin{gather*}
\nabla^{\eta}\left[w_{2}(\alpha x+\beta)\right]=w_{2+\eta}(\alpha x+\beta)+q \eta \alpha w_{1+\eta}  \tag{24}\\
\nabla^{\eta}\left[w_{1}(\gamma x+v \alpha+\delta)\right]=w_{1+\eta}(\gamma x+v \alpha+\delta)+q \eta \gamma w_{\eta} \tag{25}
\end{gather*}
$$

where $q$ is a shift operator which is defined by $w(t-1)=q w(t)$. By substituting (24), (25) into the (23), we have

$$
\begin{equation*}
w_{2+\eta}(\alpha x+\beta)+[(\eta q+v) \alpha+\gamma x+\delta] w_{1+\eta}+(q \eta \gamma+v \gamma) w_{\eta}=\psi_{\eta} \tag{26}
\end{equation*}
$$

We choose $\eta$ such that

$$
q \eta \gamma+v \gamma=0, \quad \eta=-q^{-1} v
$$

Then we obtain

$$
\begin{equation*}
w_{2-q^{-1} v}(\alpha x+\beta)+w_{1-q^{-1} v}(\gamma x+\delta)=\psi_{-q^{-1} v} \tag{27}
\end{equation*}
$$

from (26).
Therefore, setting

$$
\begin{equation*}
w_{1-q^{-1} v}=u \quad\left(w=u_{-1+q^{-1} v}\right) \tag{28}
\end{equation*}
$$

we have

$$
\begin{equation*}
u_{1}+u\left(\frac{\gamma x+\delta}{\alpha x+\beta}\right)=\psi_{-q^{-1} v}(\alpha x+\beta)^{-1} \tag{29}
\end{equation*}
$$

from (27). A particular solution of a first order ordinary differential Equation (29):

$$
\begin{equation*}
u=\left[\psi_{-q^{-1} v}(\alpha x+\beta)^{\left(\delta \alpha-\gamma \beta-\alpha^{2}\right) / \alpha^{2}} e^{\frac{\gamma}{\alpha} x}\right]_{-1}(\alpha x+\beta)^{(\gamma \beta-\delta \alpha) / \alpha^{2}} e^{-\frac{\gamma}{\alpha} x} \tag{30}
\end{equation*}
$$

Thus we obtain the solution (21) from (28) and (30).
(ii) Set

$$
\begin{equation*}
w=(\alpha x+\beta)^{\sigma} W(x) \tag{31}
\end{equation*}
$$

The first and second derivations of (31) are acquired as follows:

$$
\begin{gather*}
w_{1}=\sigma(\alpha x+\beta)^{\sigma-1} \alpha W+(\alpha x+\beta)^{\sigma} W_{1}  \tag{32}\\
w_{2}=\sigma(\sigma-1)(\alpha x+\beta)^{\sigma-2} \alpha^{2} W+2 \sigma(\alpha x+\beta)^{\sigma-1} \alpha W_{1}+(\alpha x+\beta)^{\sigma} W_{2} \tag{33}
\end{gather*}
$$

Substitute (31)-(33) into (20), we have

$$
\begin{align*}
& \quad W_{2}(\alpha x+\beta)+W_{1}(2 \sigma \alpha+\gamma x+v \alpha+\delta)  \tag{34}\\
& \\
& +W\left(\frac{\alpha^{2} \sigma(\sigma-1)+\alpha \sigma(\gamma x+v \alpha+\delta)}{\alpha x+\beta}+v \gamma\right) \\
& = \\
& \psi(\alpha x+\beta)^{-\sigma} .
\end{align*}
$$

Here, we choose $\sigma$ such that

$$
\alpha \sigma(\alpha \sigma-\alpha+\gamma x+v \alpha+\delta)=0
$$

that is $\sigma=0, \sigma=\frac{-(\gamma x+\delta)}{\alpha}-v+1$.
In the case $\sigma=0$, we have the same results as $i$.
Let $\sigma=\frac{-(\gamma x+\delta)}{\alpha}-v+1$. From (31) and (34)

$$
\begin{equation*}
w=(\alpha x+\beta)^{\frac{-(\gamma x+\delta)}{\alpha}-v+1} W \tag{35}
\end{equation*}
$$

and

$$
\begin{equation*}
W_{2}(\alpha x+\beta)+W_{1}[\alpha(2-v)-\delta-\gamma x]+v \gamma W=\psi(\alpha x+\beta)^{\frac{\gamma x+\delta}{\alpha}+v-1} \tag{36}
\end{equation*}
$$

respectively.
Applying the operator $\nabla^{\eta}$ to both members of (36), we have

$$
\begin{align*}
& W_{2+\eta}(\alpha x+\beta)+W_{1+\eta}[\alpha(2-v+\eta q)-\delta-\gamma x]+W_{\eta}(-\gamma \eta q+v \gamma)  \tag{37}\\
= & {\left[\psi(\alpha x+\beta)^{\frac{\gamma x+\delta}{\alpha}+v-1}\right]_{\eta} . }
\end{align*}
$$

Choose $\eta$ such that

$$
-\gamma \eta q+v \gamma=0, \quad \eta=q^{-1} v
$$

we have then

$$
\begin{equation*}
W_{2+q^{-1} v}(\alpha x+\beta)+W_{1+q^{-1} v}[2 \alpha-(\gamma x+\delta)]=\left[\psi(\alpha x+\beta)^{\frac{\gamma x+\delta}{\alpha}+v-1}\right]_{q^{-1} v} \tag{38}
\end{equation*}
$$

from (37).
Therefore, setting

$$
\begin{equation*}
W_{1+q^{-1} v}=\vartheta, \quad W=\vartheta_{-1-q^{-1} v} \tag{39}
\end{equation*}
$$

we have

$$
\begin{equation*}
\vartheta_{1}+\vartheta\left[\frac{2 \alpha}{\alpha x+\beta}-\frac{\gamma x+\delta}{\alpha x+\beta}\right]=\left[\psi(\alpha x+\beta)^{\frac{\gamma x+\delta}{\alpha}+v-1}\right]_{q^{-1} v}(\alpha x+\beta)^{-1} \tag{40}
\end{equation*}
$$

from (38). A particular solution of ordinary differential Equation (40) is given by

$$
\begin{equation*}
\vartheta=\left\{\left[\psi(\alpha x+\beta)^{\frac{\gamma x+\delta}{\alpha}+v-1}\right]_{q^{-1} v}(\alpha x+\beta)^{\frac{-\delta \alpha+\gamma \beta}{\alpha^{2}}+1} e^{-\frac{\gamma}{\alpha} x}\right\}_{-1}(\alpha x+\beta)^{\frac{\delta \alpha-\gamma \beta}{\alpha^{2}}-2} e^{\frac{\gamma}{\alpha} x} \tag{41}
\end{equation*}
$$

Thus we obtain the solution (22) from (35), (39) and (41).
Furthermore, we can prove for the homogen part such that the homogeneous linear ordinary differential equation

$$
\begin{equation*}
w_{2}(\alpha x+\beta)+w_{1}(\gamma x+v \alpha+\delta)+v \gamma w(x)=0, x \neq-\frac{\beta}{\alpha}, \alpha \neq 0, v \in \mathbb{R} \tag{42}
\end{equation*}
$$

has solutions of the forms

$$
\begin{gather*}
w^{\mathrm{I}}(x)=h\left[(\alpha x+\beta)^{(\gamma \beta-\delta \alpha) / \alpha^{2}} e^{-\frac{\gamma}{\alpha} x}\right]_{-1+q^{-1} v^{\prime}}  \tag{43}\\
w^{\mathrm{II}}(x)=h(\alpha x+\beta)^{\frac{-(\gamma x+\delta)}{\alpha}-v+1}\left[(\alpha x+\beta)^{\frac{\delta \alpha-\gamma \beta}{\alpha^{2}}-2} e^{\frac{\gamma}{\alpha} x}\right]_{-1-q^{-1} v} \tag{44}
\end{gather*}
$$

where $h$ is an arbitrary constant.
Now, in Theorem 1, we further set

$$
\begin{equation*}
\alpha=1, \beta=\ell, \gamma=2 \kappa+a, \delta=2 \kappa \ell-v, v=\frac{\kappa^{2} \ell+\tau a+d}{2 \kappa+a} \tag{45}
\end{equation*}
$$

and let

$$
\psi(x) \rightarrow x^{1-\tau} e^{-\kappa x} \psi(x)
$$

We thus find that the nonhomogeneous differential Equation (19) has a particular solution given by

$$
\begin{align*}
w^{\mathrm{I}}(x)= & \left(\left\{\left[x^{1-\tau} e^{-\kappa x} \psi(x)\right]_{-q^{-1} v}(x+\ell)^{-v-a \ell-1} e^{(2 \kappa+a) x}\right\}_{-1}\right. \\
& \left.(x+\ell)^{v+a \ell} e^{-(2 \kappa+a) x}\right)_{-1+q^{-1} v} \tag{46}
\end{align*}
$$

$$
\begin{align*}
w^{\mathrm{II}}(x)= & (x+\ell)^{-(2 \kappa+a) x-2 \kappa \ell+1} \\
& \times\left(\left\{\left[x^{1-\tau} e^{-\kappa x} \psi(x)(x+\ell)^{(2 \kappa+a) x+2 \kappa \ell-1}\right]_{q^{-1} v}(x+\ell)^{v+a \ell+1} e^{-(2 \kappa+a) x}\right\}_{-1}\right. \\
& \left.\times(x+\ell)^{-v-a \ell-2} e^{(2 \kappa+a) x}\right)_{-1-q^{-1} v} \tag{47}
\end{align*}
$$

and that the corresponding homogeneous linear differential equation

$$
\begin{equation*}
(x+\ell) \frac{d^{2} w}{d x^{2}}+[2 \kappa \ell+(2 \kappa+a) x] \frac{d w}{d x}+\left(\kappa^{2} \ell+\tau a+d\right) w(x)=0 \tag{48}
\end{equation*}
$$

has solutions of the forms

$$
\begin{gather*}
w^{\mathrm{I}}(x)=h\left[(x+\ell)^{v+a \ell} e^{-(2 \kappa+a) x}\right]_{-1+q^{-1} v}  \tag{49}\\
w^{\mathrm{II}}(x)=h(x+\ell)^{-(2 \kappa+a) x-2 \kappa \ell+1}\left[(x+\ell)^{-v-a \ell-2} e^{(2 \kappa+a) x}\right]_{-1-q^{-1} v} \tag{50}
\end{gather*}
$$

where $h$ is an arbitrary constant.
Therefore, the linear differential Equation (11), has a particular solution in the following forms

$$
\begin{align*}
y^{\mathrm{I}}(x)= & x^{\tau} e^{\kappa x} w(x) \\
= & x^{\tau} e^{\kappa x}\left(\left\{\left[x^{1-\tau} e^{-\kappa x} \psi(x)\right]_{-q^{-1} v}(x+\ell)^{-v-a \ell-1} e^{(2 \kappa+a) x}\right\}_{-1}\right. \\
& \left.(x+\ell)^{v+a \ell} e^{-(2 \kappa+a) x}\right)_{-1+q^{-1} v} \quad x \in \mathbb{C} \backslash\{0,-\ell\}, v \in \mathbb{R} \tag{51}
\end{align*}
$$

and

$$
\begin{align*}
y^{\mathrm{II}}(x)= & x^{\tau} e^{\kappa x}(x+\ell)^{-(2 \kappa+a) x-2 \kappa \ell+1} \\
& \times\left(\left\{\left[x^{1-\tau} e^{-\kappa x} \psi(x)(x+\ell)^{(2 \kappa+a) x+2 \kappa \ell-1}\right]_{q^{-1} v}(x+\ell)^{v+a \ell+1} e^{-(2 \kappa+a) x}\right\}_{-1}\right. \\
& \left.\times(x+\ell)^{-v-a \ell-2} e^{(2 \kappa+a) x}\right)_{-1-q^{-1} v} \tag{52}
\end{align*}
$$

and that the corresponding homogeneous linear differential equation

$$
\begin{equation*}
\left(1+\frac{\ell}{x}\right) \frac{d^{2} y}{d x^{2}}+\left[a+\frac{b}{x}\left(1+\frac{\ell}{x}\right)\right] \frac{d y}{d x}+\left[c+\frac{d}{x}+\frac{\varepsilon}{x^{2}}\left(1+\frac{\ell}{x}\right)\right] y(x)=0 \tag{53}
\end{equation*}
$$

has solutions given by

$$
\begin{gather*}
y^{\mathrm{I}}(x)=h x^{\tau} e^{\kappa x}\left[(x+\ell)^{v+a \ell} e^{-(2 \kappa+a) x}\right]_{-1+q^{-1} v^{\prime}}  \tag{54}\\
y^{\mathrm{II}}(x)=h x^{\tau} e^{\kappa x}(x+\ell)^{-(2 \kappa+a) x-2 \kappa \ell+1}\left[(x+\ell)^{-v-a \ell-2} e^{(2 \kappa+a) x}\right]_{-1-q^{-1} v} \tag{55}
\end{gather*}
$$

where $h \in \mathbb{R}$, the parameters $\tau, \kappa$ and $v$ are given by (17), (18) and (45).
Remark 1. First of all, when $\ell=0$, the differential Equation (11) reduces to the following version of the Tricomi equation:

$$
\frac{d^{2} y}{d x^{2}}+\left(a+\frac{b}{x}\right) \frac{d y}{d x}+\left(c+\frac{d}{x}+\frac{\varepsilon}{x^{2}}\right) y=\psi(x)
$$

By setting

$$
\ell=0, \quad a=b=0, \quad c=k^{2}, \quad d=n, \quad \varepsilon=\frac{1}{4}-m^{2}
$$

in the Equation (11), we readily obtain the following Hydrogen atom equation:

$$
\frac{d^{2} y}{d x^{2}}+\left(k^{2}+\frac{n}{x}+\frac{\frac{1}{4}-m^{2}}{x^{2}}\right) y=\psi(x)
$$

Example 1. In the case $\alpha=1, \beta=\gamma=v=0, \delta=2$ and $\psi(x)=x$, we have

$$
\begin{equation*}
w(x)+\frac{2}{x} w_{1}=1 \quad(x \neq 0) \tag{56}
\end{equation*}
$$

from (20). Solution of Equation (56) is obtained as

$$
\begin{align*}
& w(x)=\left\{\left[x^{2}\right]_{-1} x^{-2}\right\} \\
= & \left\{\frac{x^{3}}{3} x^{-2}\right\}  \tag{57}\\
= & \frac{1}{6} x^{2}
\end{align*}
$$

by using (21). The function obtained in (57) provide the Equation (56).

## 4. Conclusions

In this article, we use the discrete fractional operator for the homogeneous and non-homogeneous non-Fuchsian differential equations. This solution of the equation has not been obtained before by using $\nabla$ operator. We can obtain particular solutions of the same type linear singular ordinary and partial differential equations by using the discrete fractional nabla operator in future works.

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